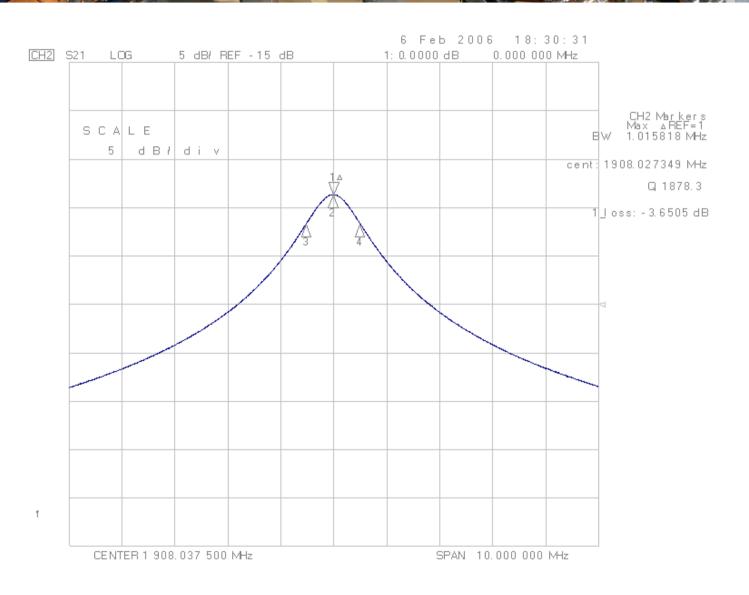
Cavity transmission measurement

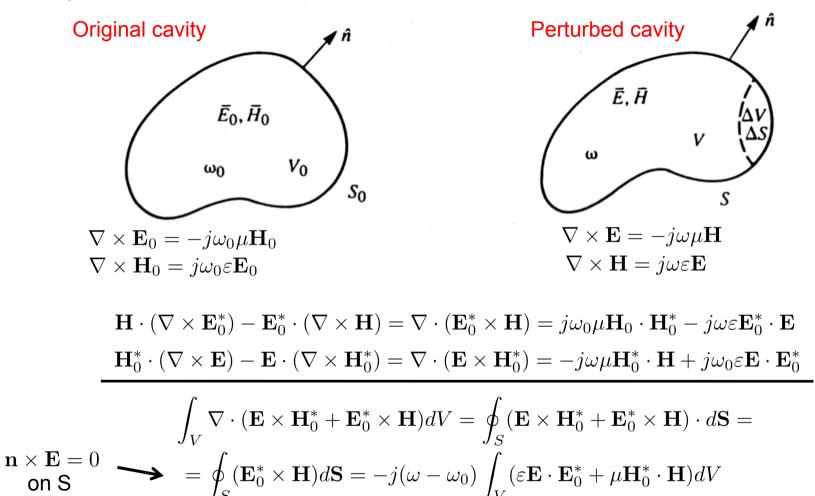


Slater theorem and applications

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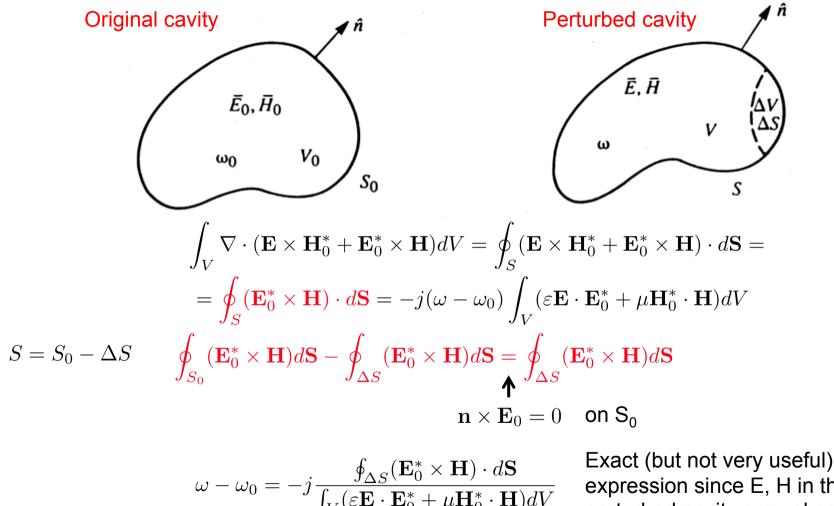
Inserting small metallic objects into a cavity or slightly deforming the shape of the cavity can be treated by the perturbation technique (**Slater theorem**)



See H. Henke, RF Engineering, CAS 2000



Inserting small metallic objects into a cavity or slightly deforming the shape of the cavity can be treated by the perturbation technique (Slater theorem)



expression since E, H in the perturbed cavity are unknown.

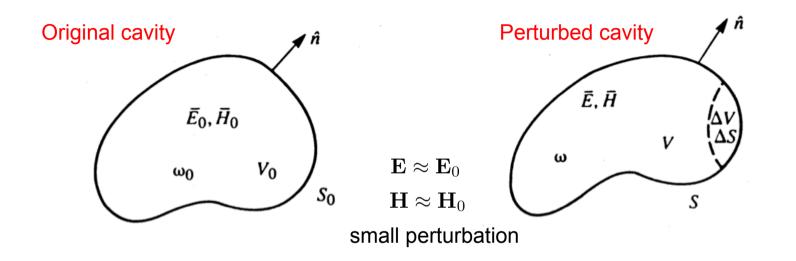
See H. Henke, RF Engineering, CAS 2000

Cavity shape perturbation 1 ĥ Perturbed cavity **Original cavity** $\overline{E}, \overline{H}$ $\overline{E}_0, \overline{H}_0$ _{V0}) $\mathbf{E} pprox \mathbf{E}_0$ Sn $\mathbf{H} \approx \mathbf{H}_0$ S small perturbation → Poyting Theor. $\oint_{A,C} (\mathbf{E}_0^* \times \mathbf{H}) \cdot d\mathbf{S} \approx \oint_{A,C} (\mathbf{E}_0^* \times \mathbf{H}_0) \cdot d\mathbf{S} = -j\omega_0 \int_{A,V} (\varepsilon |\mathbf{E}_0|^2 - \mu |\mathbf{H}_0|^2) dV$ $\omega - \omega_0 = -j \frac{\oint_{\Delta S} (\mathbf{E}_0^* \times \mathbf{H}) \cdot d\mathbf{S}}{\int_{\mathbf{U}} (\varepsilon \mathbf{E} \cdot \mathbf{E}_0^* + \mu \mathbf{H}_0^* \cdot \mathbf{H}) dV}$ It is essentially the energy stored in the cavity and it will not change Change in stored electric/ much with the perturbation. magnetic energy $\frac{\omega - \omega_0}{\omega_0} \approx \frac{\int_{\Delta V} (\mu |\mathbf{H}_0|^2 - \varepsilon |\mathbf{E}_0|^2) dV}{\int_{V_0} (\varepsilon |\mathbf{E}_0|^2 + \mu |\mathbf{H}_0|^2) dV} = \frac{\Delta W_m - \Delta W_e}{W_m + W_e}$

Total energy stored

See H. Henke, RF Engineering, CAS 2000



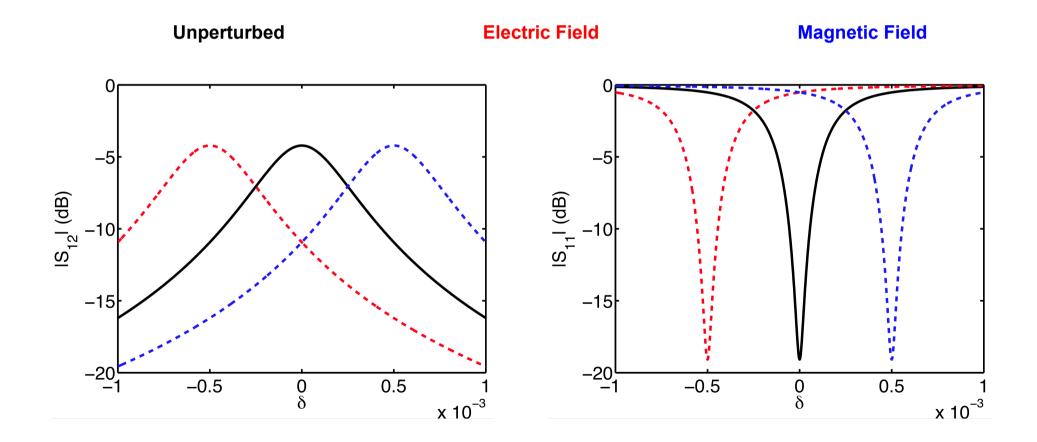


$$\frac{\Delta\omega}{\omega_0} \approx \frac{\int_{\Delta V} (\mu |\mathbf{H}_0|^2 - \varepsilon |\mathbf{E}_0|^2) dV}{\int_{V_0} (\varepsilon |\mathbf{E}_0|^2 + \mu |\mathbf{H}_0|^2) dV} = \frac{1}{W_{tot}} \int_{\Delta V} \mu |\mathbf{H}_0|^2 dV - \frac{1}{W_{tot}} \int_{\Delta V} \varepsilon |\mathbf{E}_0|^2 dV$$

The frequency shift depends on the kind and the amplitude of the unperturbed cavity field.

Applications: tuning of a cavity (S parameters)

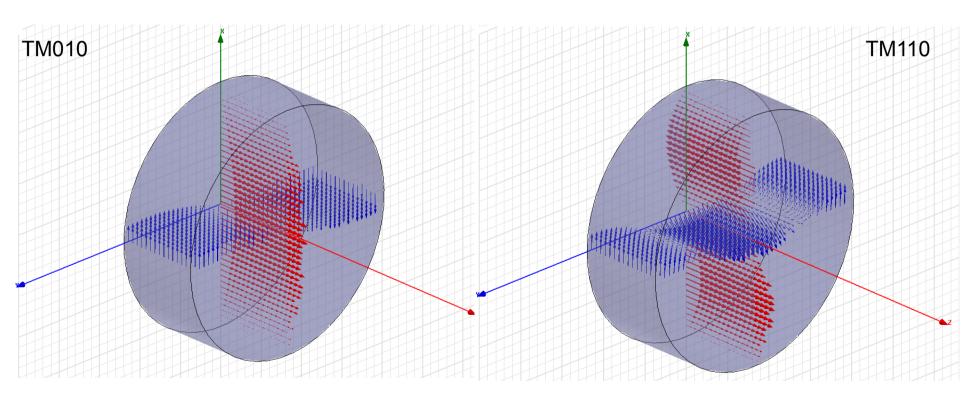
$$\frac{\Delta\omega}{\omega_0} \approx \frac{1}{W_{tot}} \int_{\Delta V} \mu |\mathbf{H}_0|^2 dV - \frac{1}{W_{tot}} \int_{\Delta V} \varepsilon |\mathbf{E}_0|^2 dV$$



Applications: tuning of different cavity modes

The same tuners affect different modes in different ways ...

$$\frac{\Delta\omega}{\omega_0} \approx \frac{1}{W_{tot}} \int_{\Delta V} \mu |\mathbf{H}_0|^2 dV - \frac{1}{W_{tot}} \int_{\Delta V} \varepsilon |\mathbf{E}_0|^2 dV$$



Task: for a given cavity with two tuners, you can tune the resonant frequency two modes simultaneously.

Courtesy of L. Ficcadenti

Applications: bead pull measurement

Introducing a small metallic object into the interior of the cavity should perturb the frequency in a similar way by an amount depending upon the local fields, and thus we could use the **frequency shift to measure the field strength** at an interior point.

$$\frac{\Delta\omega}{\omega_{0}} \approx \frac{1}{W_{tot}} \int_{\Delta V} \mu |\mathbf{H}_{0}|^{2} dV - \frac{1}{W_{tot}} \int_{\Delta V} \varepsilon |\mathbf{E}_{0}|^{2} dV$$
Measurements
$$\frac{\Delta\omega}{\omega_{0}} \approx \left(k_{\parallel}^{H} \mu \frac{|H_{z}|^{2}}{W_{tot}} + k_{\perp}^{H} \mu \frac{|\mathbf{H}_{\perp}|^{2}}{W_{tot}} \right) - \left(k_{\parallel}^{E} \varepsilon \frac{|E_{z}|^{2}}{W_{tot}} + k_{\perp}^{E} \varepsilon \frac{|\mathbf{E}_{\perp}|^{2}}{W_{tot}} \right)$$
Theory and/or calibration in known cavities
$$\frac{|E_{z}|^{2}}{W_{tot}} \approx -\frac{1}{k_{\perp}^{E} \varepsilon} \frac{\Delta\omega}{\omega_{0}} \qquad \text{Measurements}$$
Advancements
$$\frac{|E_{z}|^{2}}{W_{tot}} \approx -\frac{1}{k_{\perp}^{E} \varepsilon} \frac{\Delta\omega}{\omega_{0}} \qquad \text{Measurements}$$

Applications: bead pull measurement

Introducing a small metallic object into the interior of the cavity should perturb the frequency in a similar way by an amount depending upon the local fields, and thus we could use the **frequency shift to measure the field strength** at an interior point.

$$\frac{\Delta\omega}{\omega_{0}} \approx \frac{1}{W_{tot}} \int_{\Delta V} \mu |\mathbf{H}_{0}|^{2} dV - \frac{1}{W_{tot}} \int_{\Delta V} \varepsilon |\mathbf{E}_{0}|^{2} dV$$
Measurements
$$\frac{\Delta\omega}{\omega_{0}} \approx \left(k_{\parallel}^{H} \mu \frac{|H_{z}|^{2}}{W_{tot}} + k_{\perp}^{H} \mu \frac{|\mathbf{H}_{\perp}|^{2}}{W_{tot}} \right) - \left(k_{\parallel}^{E} \varepsilon \frac{|E_{z}|^{2}}{W_{tot}} + k_{\perp}^{E} \varepsilon \frac{|\mathbf{E}_{\perp}|^{2}}{W_{tot}} \right)$$
Theory and/or calibration in known cavities

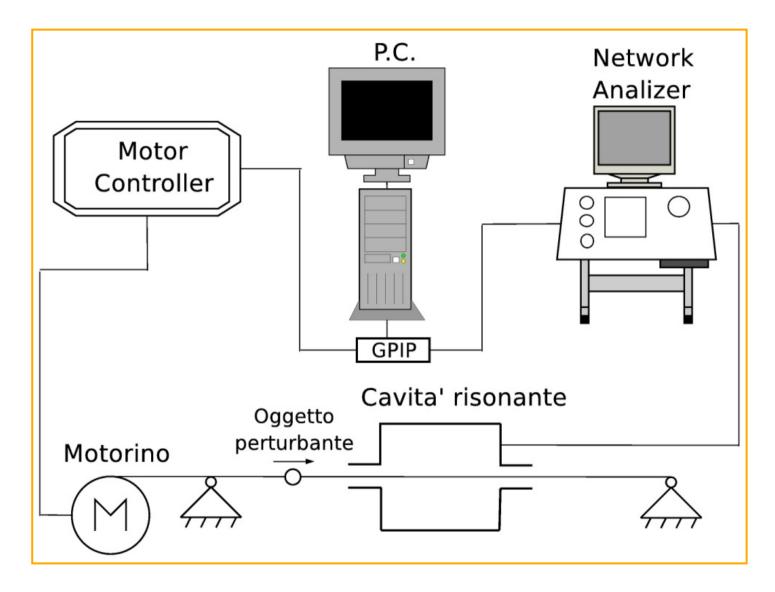
In general

Metal objects affect E and H field

Dielectric objects affect E field

It is possible to measure H field with two measurements

Automatic field measurement (bead pull

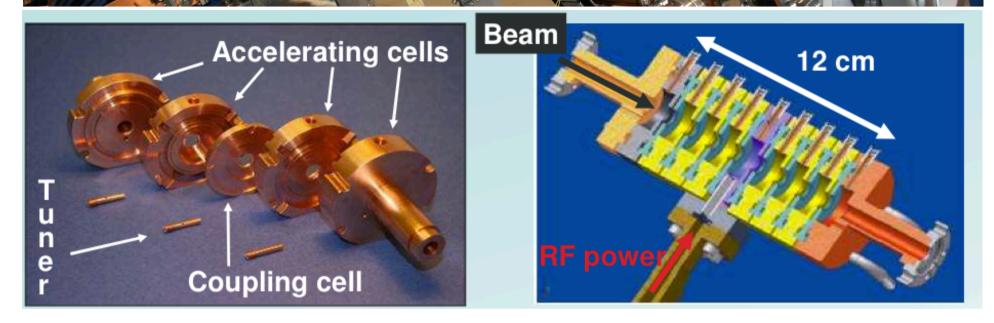


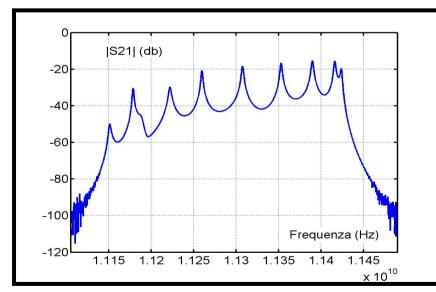
Courtesy of L. Ficcadenti

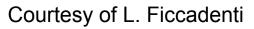
Examples

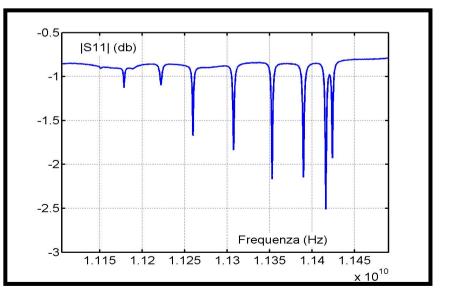
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Periodic cavity: S parameters

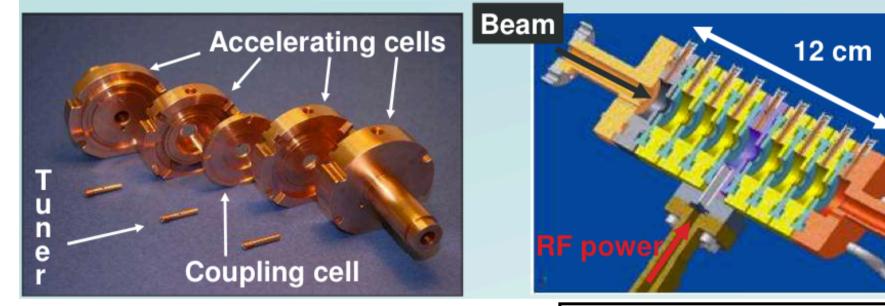


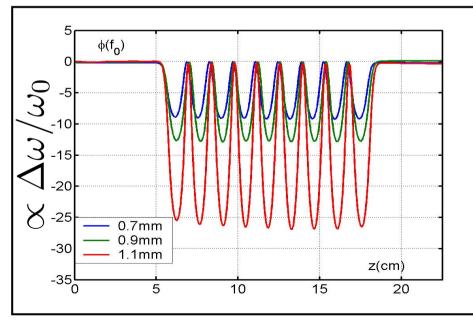


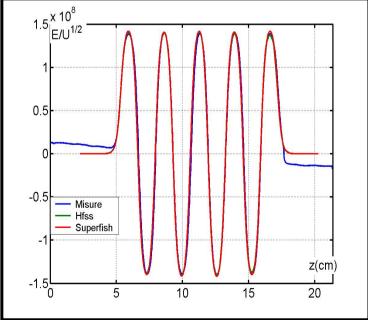




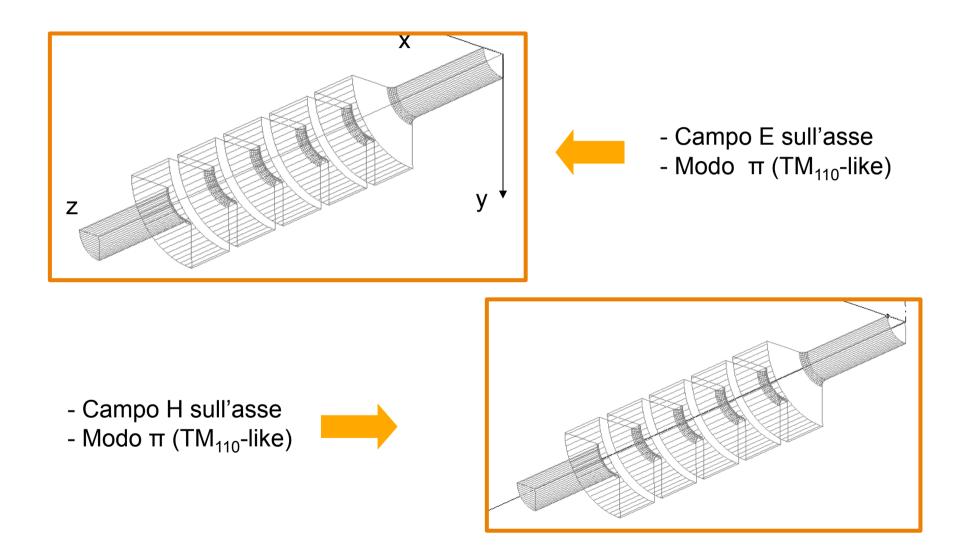
Periodic cavities: field on axis











Courtesy of L. Ficcadenti

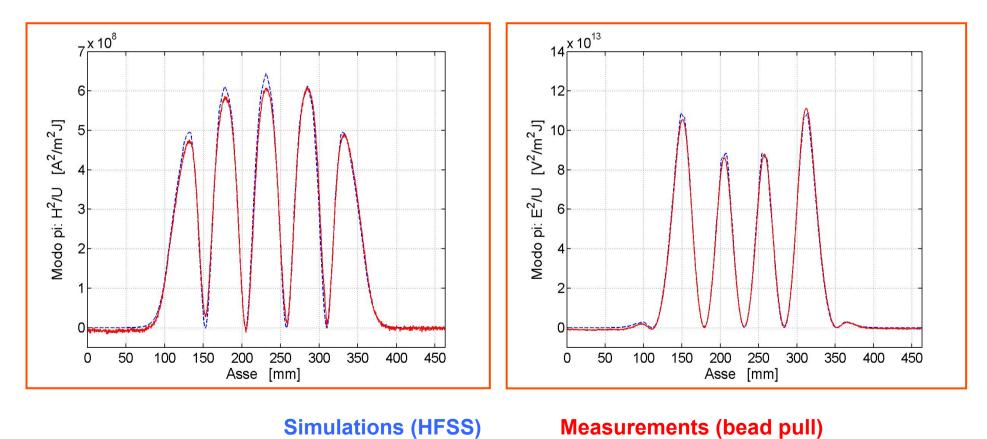






Courtesy of L. Ficcadenti

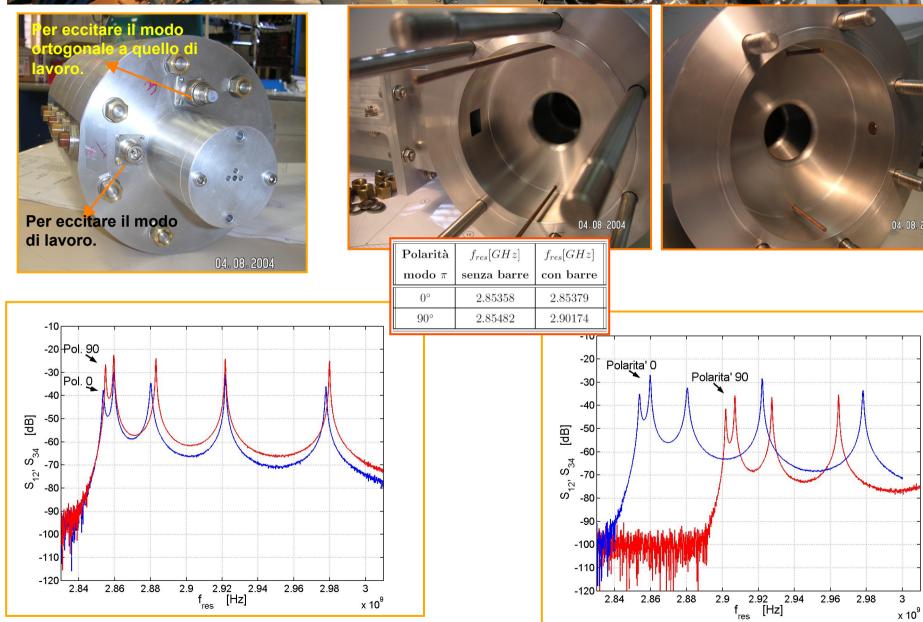




Measurements (bead pull)

Courtesy of L. Ficcadenti

RF deflector field polarities



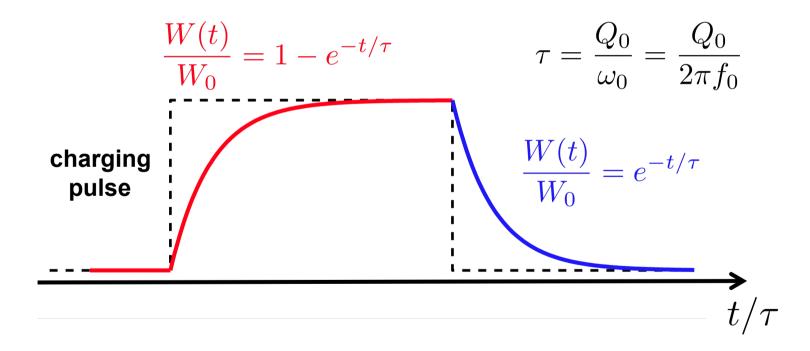
Courtesy of L. Ficcadenti

Cavity transient

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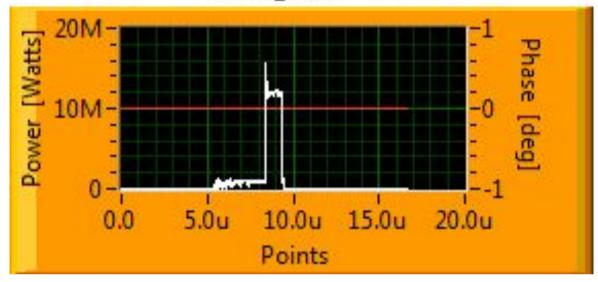


$$-dW = P_d dt = \frac{\omega_0 W}{Q_0} dt$$





DGUN_FRW



GunProbe

